

Lecture 3

Divergence free Kronecker webs in 4D.
Relation with heavenly PDEs and vacuum self-dual
Einstein metrics in neutral signature.

Mason-Newman formalism (1989): Thm. Let ω be a volume form on M^4 , $X_1(\lambda) = V_1 + \lambda V_2$, $X_2(\lambda) = V_3 + \lambda V_4$ parameter depending vector fields on M such that

(1) $[X_1(\lambda), X_2(\lambda)] = 0 \iff [V_1, V_2] = 0, [V_2, V_4] = 0, [V_1, V_4] + [V_2, V_3] = 0$

(2) $\text{div}_\omega X_i(\lambda) = 0 \iff \text{div}_\omega V_i = 0$.

Then $g = \alpha (V^1 \otimes V^4 - V^2 \otimes V^3)$, where $\alpha = \omega(V_1, V_2, V_3, V_4)$, (V^i) is the coframe dual to the frame (V_i) , is self-dual (of signature (2,2)) and satisfies vacuum Einstein equations. This correspondence is one-to-one.

Discussion on self-duality etc.: $(M^4, g) \rightsquigarrow * : \Lambda^2 T^*M \rightarrow \Lambda^2 T^*M$

Signatures $(++++), (++)-- \implies *^2 = 1$ Hodge star
 $\Lambda^2 T^*M = \Lambda^+ T^*M \oplus \Lambda^- T^*M$
 Signature $(+---) \implies *^2 = -1$ (no self-duality)

$*^2 = 1$ case : $R_{abcd} = R_{[c]a[b]d}$ $\implies R : \Lambda^+ T^*M \rightarrow \Lambda^+ T^*M$

$$R = \begin{bmatrix} C_+ - \frac{R}{12} & \psi \\ \psi & C_- - \frac{R}{12} \end{bmatrix}$$

ψ - traceless Ricci tensor

C_\pm - (anti) self-dual Weyl tensor

R - Ricci scalar

The metric is selfdual vacuum Einstein if

$$R = \begin{bmatrix} C_+ & 0 \\ 0 & 0 \end{bmatrix}.$$

Plebanski heavenly equations (1975):

I Plebanski equation $f_{13}f_{24} - f_{14}f_{23} = 1$ ($f_{ij} = \frac{\partial^2 \mathcal{F}}{\partial x_i \partial x_j}$)

$$X_1(\lambda) = -\partial_3 + \lambda(f_{13}\partial_2 - f_{23}\partial_1)$$

$$X_2(\lambda) = -\partial_4 + \lambda(f_{14}\partial_2 - f_{24}\partial_1)$$

$$\omega = dx_1 \wedge dx_2 \wedge dx_3 \wedge dx_4$$

$$g = 2(f_{13} dx_1 dx_3 + f_{14} dx_1 dx_4 + f_{23} dx_2 dx_3 + f_{24} dx_2 dx_4)$$

Other equations leading to Mason-Newman vector fields:

Plebanski II equation (1975) $f_{13} + f_{24} + f_{11} f_{22} - f_{12}^2 = 0$

Husain-Park equation (1992-94) $f_{34} + f_{23} f_{14} - f_{13} f_{24} = 0$

Schief or general heavenly equation (1996)

$$(\lambda_1 - \lambda_2)(\lambda_3 - \lambda_4) f_{12} f_{34} + (\lambda_1 - \lambda_3)(\lambda_4 - \lambda_2) f_{13} f_{24} + (\lambda_1 - \lambda_4)(\lambda_2 - \lambda_3) f_{14} f_{23} = 0$$

Interpretation from the point of view of Kronecker webs:

Mason-Newman vector fields $X_1(\lambda), X_2(\lambda)$ generate a Kronecker web $\{F_\lambda\}$ of codim 2 on M^4 (two Kronecker blocks of the type $\begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix}$) with the additional condition $\text{div}_\omega X_i(\lambda) = 0$.

Call such Kronecker webs divergence free.

Remarks 0. It is easy to show that a Kronecker web $\{F_\lambda\}$ of codim 2 on M^4 is divergence free $\Leftrightarrow T F_\lambda = \ker \beta_\lambda$, where $\beta^\lambda = \beta_0 + \lambda \beta_1 + \lambda^2 \beta_2$ is a 2-form such that $\beta^\lambda \wedge \beta^\lambda = 0, d\beta^\lambda = 0$.

Gindikin 1982

1. Fixing any $\lambda_1, \lambda_2, \lambda_3$ pairwise distinct one gets a classical 3-web $\{F_{\lambda_1}, F_{\lambda_2}, F_{\lambda_3}\}$.
2. One can calculate Chern connection for this 3-web, it will be torsionless. Moreover the divergence free condition implies that the Ricci tensor of this connection vanishes.
3. One can choose the parameters $\lambda_1, \lambda_2, \lambda_3$ in such a way that the Chern connection will coincide with the Levi-Civita connection of the corresponding metric g .

Konopelchenko-Schief-Szerczewski observations (2021):

Let $X_1(\lambda), X_2(\lambda)$ be a Mason-Newman vector fields

1. Fix four pairwise distinct parameters $\lambda_1, \lambda_2, \lambda_3, \lambda_4$ and use common eigenfunctions φ_i for the vector fields $X_1(x_i), X_2(x_i)$ as coordinates. Then the PNO corresponding to the Kronecker web $\{F_\lambda\}$ is given by $N: T\mathbb{F}_\infty \rightarrow TM, N = \tilde{N} / T\mathbb{F}_\infty$, where $\tilde{N}: TM \rightarrow TM$ is the Nijenhuis operator $\frac{\partial}{\partial \varphi_i} \mapsto \lambda_i \frac{\partial}{\partial \varphi_i}$ (no summation).

2. There exists a "potential", i.e. a function f such that it satisfies Schief equation and the initial Maschke-Newman vector fields are given by

$$X_1(\lambda) = \frac{f_{24}(\lambda_2 - \lambda_4)}{f_{12}(\lambda_1 - \lambda_2)} (\lambda_1 - \lambda) \partial_1 - \frac{f_{14}(\lambda_1 - \lambda_4)}{f_{12}(\lambda_1 - \lambda_2)} (\lambda_2 - \lambda) \partial_2 + (\lambda_4 - \lambda) \partial_4, \quad X_2(\lambda) = \frac{f_{23}(\lambda_2 - \lambda_3)}{f_{12}(\lambda_1 - \lambda_2)} (\lambda_1 - \lambda) \partial_1 - \frac{f_{13}(\lambda_1 - \lambda_3)}{f_{12}(\lambda_1 - \lambda_2)} (\lambda_2 - \lambda) \partial_2 + (\lambda_3 - \lambda) \partial_3$$

and ω is given by $\omega = f_{12} d\varphi_1 \wedge d\varphi_2 \wedge d\varphi_3 \wedge d\varphi_4$

3. If a function f satisfies dispersionless Hirota system of PDEs then it satisfies the Schief equation.

Veronese webs in 4D \iff divergence free Kronecker webs:

$$\{V_\lambda\}, TV_\lambda = \ker \alpha^\lambda$$

$$\alpha^\lambda = \lambda_0 + \lambda \lambda_1 + \lambda^2 \lambda_2 + \lambda^3 \lambda_3$$

or in the Zakharov approach

$$\alpha^\lambda = (\lambda - \lambda_1)(\lambda - \lambda_2)(\lambda - \lambda_3)(\lambda - \lambda_4) \times \underbrace{\left(\frac{f_1 dx_1}{\lambda - \lambda_1} + \frac{f_2 dx_2}{\lambda - \lambda_2} + \frac{f_3 dx_3}{\lambda - \lambda_3} + \frac{f_4 dx_4}{\lambda - \lambda_4} \right)}_{((\tilde{N} - \lambda I)^{-1})^t d f}$$

$$\{F_\lambda\}, TF_\lambda = \ker \beta^\lambda$$

$$\beta^\lambda = \beta_0 + \lambda \beta_1 + \lambda^2 \beta_2$$

in the K-S-S approach

$$\beta^\lambda = (\lambda - \lambda_1)(\lambda - \lambda_2)(\lambda - \lambda_3)(\lambda - \lambda_4) \times \underbrace{\sum_{i < j} \frac{f_{ij} dx_i \wedge dx_j}{(\lambda - \lambda_i)(\lambda - \lambda_j)}}_{d((\tilde{N} - \lambda I)^{-1})^t d f}$$

$$\tilde{N}: T\mathbb{R}^4 \rightarrow T\mathbb{R}^4 \quad \frac{\partial}{\partial x_i} \mapsto \lambda_i \frac{\partial}{\partial x_i}$$

$$(*) \quad \alpha^\lambda \wedge d\alpha^\lambda = 0 \iff \text{Hirota system} \xrightarrow{d(*)} \beta^\lambda \wedge \beta^\lambda = 0 \iff \text{Schief eq.}$$

$$\beta^\lambda = d\alpha^\lambda$$

$$(d\alpha_3 = 0 \text{ since } \alpha_3 = d\ell)$$

Moreover $\alpha^\lambda \wedge \beta^\lambda = 0$ implies that

$$TV_\lambda \supset TF_\lambda$$

Remark In fact from a given Veronese web we obtain a 1-parametric family of divergence free Kromer webs:

$$\alpha^\lambda \wedge d\alpha^\lambda = 0 \implies \forall \lambda = \lambda_0 \exists \text{ integrating multiplier } g^{\lambda_0} \text{ such that}$$

$$d(g^{\lambda_0} \alpha^{\lambda_0}) = 0$$

$$\Downarrow$$

$$\beta^\lambda = d(g^{\lambda_0} \alpha^\lambda) \text{ is of degree 2 in } \lambda.$$

The forms β^{λ_0} differ essentially for different λ_0 : $d(g \cdot \alpha^0) = dg \alpha^0 + g d\alpha^0$

New integrable "heavenly" PDEs (P-Szeereszewski, 2022):

Consider different extensions of the PND $\tilde{N}|_{TF_0}: TF_0 \rightarrow TM^4$ to globally defined Nijenhuis operators.

Each normal form gives its own integrable PDE.

Open questions:

1. Give a concise characterization of vacuum self dual Einstein metrics coming from Veronese webs
2. Study transformation of canonical connections

- on Veronese web to canonical connections
 - or divergence free Kronecker webs
3. Study particular solutions